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COMMENT

## On self-avoiding walks in critical dimensions

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**Abstract.** A new interpretation of the numerical data for self-avoiding walks in critical dimensions is suggested on the basis of a different renormalization scheme for the random walk with a long-term correlation.

In a recent paper [1] Grassberger *et al* presented very rich numerical data on the self-avoiding walks in critical dimensions  $d = 4$ . They clearly demonstrated the logarithmic dependence predicted by the standard renormalization theory. However, the critical exponent of this dependence,  $\alpha = 0.31$ , found by the fitting of numerical data to the asymptotic relation ( $\ln N \rightarrow \infty$ )

$$\langle R_N^2 \rangle = RN \left[ \ln \left( \frac{N}{A} \right) \right]^\alpha \quad (1)$$

differs substantially from the theoretical value  $\alpha = \frac{1}{4}$ . The authors [1] resolve this apparent contradiction by including the first correction term of the renormalization theory

$$\langle R_N^2 \rangle = rN \left[ \ln \left( \frac{N}{a} \right) \right]^{1/4} \left[ 1 - \frac{17 \ln(4 \ln(N/a)) + 31}{64 \ln(N/a)} \right] \quad (2)$$

with fitting parameters  $r = 1.331$  and  $a = 0.1237$  in the range  $N = 20$ –4000.

Additional numerical data up to  $N = 10^7$  are mentioned (but not given) in [1], and the local exponent is said to decrease down to  $\alpha \approx 0.285$ . Meanwhile, using three-parameter equation (1) it is possible to fit *all* the numerical data with a single value of  $\alpha = 0.298$ , and  $R = 1.10$ ,  $A = 1.22$  provided the additional data ( $N = 4000$ – $10^7$ ) match (2) with the same  $r$  and  $a$ . The fitting accuracy is  $|\delta q|/q < 1.3 \times 10^{-3}$  where  $q(N) = \langle R_N^2 \rangle$ . So, the comparison with the theory is a tricky task indeed [1].

The main purpose of this comment is to point out a different approach to the problem based upon another renormalization theory for the random walk with a long-term correlation [2, 3]. The main idea of this approach is in that the self-avoiding of a trajectory with a *finite* width results in a correlation proportional to the ratio of the trajectory length  $N$  to the volume occupied by the diffusion [3]:

$$C(N) = b \frac{N}{q^{d/2}} = \frac{d^2 q(N)}{2dN^2} \quad (3)$$

where  $b$  stands for some numerical factor and where the latter equality is the standard correlation/dispersion relation if the correlation integral (diffusion rate) diverges.

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The power-law solution to equation (3)

$$q(N) = RN^{2\nu} \quad \nu = \frac{3}{2+d} \quad R^{1+d/2} = \frac{b(2+d)^2}{3(4-d)} \quad (4)$$

leads to the well known Flory formula for  $\nu$  which is valid for  $d < 4$ . In the critical case  $d = 4$  the asymptotic solution has the form (1) with  $R = (6b)^{1/3}$  and  $\alpha = \frac{1}{3}$ . That the latter value is well in agreement with the numerical data is the main result of this comment. The remaining parameter  $A$  is determined by a non-asymptotic correction term similar to that in (2).

## References

- [1] Grassberger P, Hegger R and Schäfer L 1994 *J. Phys. A: Math. Gen.* 27 7265
- [2] Chirikov B V and Shepelyansky D L 1984 *Physica* 13D 395
- [3] Chirikov B V 1991 *Chaos Solitons Fractals* 1 79